#### Chapter 1. Heat Equation

#### Section 1.2

- 1.2.9 (d) Circular cross section means that P = 2πr, A = πr<sup>2</sup>, and thus P/A = 2/r, where r is the radius. Also  $\gamma = 0$ .
- 1.2.9 (e) u(x, t) = u(t) implies that

$$c\rho \frac{du}{dt} = -\frac{2h}{r}u$$
.

The solution of this first-order linear differential equation with constant coefficients, which satisfies the initial condition  $u(0) = u_0$ , is

$$u(t) = u_0 \exp \left[-\frac{2h}{c\rho r}t\right].$$

#### Section 1.3

1.3.2 ∂u/∂x is continuous if K<sub>0</sub>(x<sub>0</sub>−) = K<sub>0</sub>(x<sub>0</sub>+), that is, if the conductivity is continuous.

#### Section 1.4

- 1.4.1 (a) Equilibrium satisfies (1.4.14), d<sup>2</sup>u/dx<sup>2</sup> = 0, whose general solution is (1.4.17), u = c<sub>1</sub> + c<sub>2</sub>x. The boundary condition u(0) = 0 implies  $c_1 = 0$  and u(L) = T implies  $c_2 = T/L$  so that u = Tx/L.
- 1.4.1 (d) Equilibrium satisfies (1.4.14), d<sup>2</sup>u/dx<sup>2</sup> = 0, whose general solution (1.4.17), u = c<sub>1</sub> + c<sub>2</sub>x. From the boundary conditions, u(0) = T yields  $T = c_1$  and  $du/dx(L) = \alpha$  yields  $\alpha = c_2$ . Thus  $u = T + \alpha x$ .
- 1.4.1 (f) In equilibrium, (1.2.9) becomes  $d^2u/dx^2 = -Q/K_0 = -x^2$ , whose general solution (by integrating twice) is  $u = -x^4/12 + c_1 + c_2x$ . The boundary condition u(0) = T yields  $c_1 = T$ , while du/dx(L) = 0yields  $c_2 = L^3/3$ . Thus  $u = -x^4/12 + L^3x/3 + T$ .
- 1.4.1 (h) Equilibrium satisfies  $d^2u/dx^2 = 0$ . One integration yields  $du/dx = c_2$ , the second integration document is available on https://unihelp.xyz/ \*\*\* contact me if site not loaded x=0:  $c_2-(c_1-T)=0$  x=L:  $c_2=\alpha$  and thus  $c_3=T+\alpha$ .

$$x = 0$$
:  $c_2 - (c_1 - T) = 0$   
 $x = L$ :  $c_3 = \alpha$  and thus  $c_4 = T + \alpha$ .

Therefore,  $u = (T + \alpha) + \alpha x = T + \alpha(x + 1)$ .

1.4.7 (a) For equilibrium:

$$\frac{d^2u}{dx^2} = -1$$
 implies  $u = -\frac{x^2}{2} + c_1x + c_2$  and  $\frac{du}{dx} = -x + c_1$ .

From the boundary conditions  $\frac{du}{dx}(0) = 1$  and  $\frac{du}{dx}(L) = \beta$ ,  $c_1 = 1$  and  $-L + c_1 = \beta$  which is consistent only if  $\beta + L = 1$ . If  $\beta = 1 - L$ , there is an equilibrium solution ( $u = -\frac{\omega^2}{2} + x + c_2$ ). If  $\beta \neq 1 - L$ , there isn't an equilibrium solution. The difficulty is caused by the heat flow being specified at both ends and a source specified inside. An equilibrium will exist only if these three are in balance. This balance can be mathematically verified from conservation of energy;

$$\frac{d}{dt} \int_{0}^{L} c\rho u \ dx = -\frac{du}{dx}(0) + \frac{du}{dx}(L) + \int_{0}^{L} Q_{0} \ dx = -1 + \beta + L.$$

If  $\beta + L = 1$ , then the total thermal energy is constant and the initial energy = the final energy:

$$\int_{0}^{L} f(x) dx = \int_{0}^{L} \left(-\frac{x^{2}}{2} + x + c_{2}\right) dx, \text{ which determines } c_{2}.$$

If  $\beta + L \neq 1$ , then the total thermal energy is always changing in time and an equilibrium is never reached.

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